AGS APPLIED MATERIALS **XINTERFACES**

NiO Hierarchical Nanorings on SiC: Enhancing Relaxation to Tune Microwave Absorption at Elevated Temperature

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S Supporting Information

ABSTRACT: We fabricated NiO nanorings on SiC, a novel hierarchical architecture, by a facile two-step method. The dielectric properties depend on temperature and frequency in the range from 373 to 773 K and X band. The imaginary part and loss tangent increase more than four times and three times with increasing temperature, respectively. The architecture demonstrates multirelaxation and possesses high-efficient absorption. The reflection loss exceeds −40 dB and the bandwidth covers 85% of X band (approximately −20 dB). The synergistic effect between multirelaxation and conductance is beneficial to the microwave absorption. Our findings provide a novel and feasible strategy to tune microwave absorption.

KEYWORDS: NiO nanorings, hierarchical architecture, dielectric properties, multirelaxation, microwave absorption

Microwave absorber with broad frequency and high
efficiency is a long sought-after goal of researchers
because of the great potential applications reaging from military because of the great potential applications ranging from military to civil applications, such as information security, electronic countermeasures, EMI shielding, and healthcare. In recent years, plenty of new-type materials have been developed extensively and their microwave absorption is renewing rapidly. Carbon nanomaterials, including carbon nanotubes, carbon nanocoils, graphene, and their composite structures, exhibit potential applications in microwave attenuation.^{1−6} Magnetic material is another selectable microwave absorber, because of its magnetic and dielectric properties. Dielec[tr](#page-4-0)i[c](#page-4-0)−magnetic hybrid structures, are also considered to develop microwave absorption.^{7−12} The highly efficient absorbers have been intensively researched. However, optional microwave absorbers to meet t[he ha](#page-4-0)rsh requirement of thermal environments are rarely reported.

As a promoting option, SiC has been researched to improve its microwave absorption. More recently, Kuang et al. and Wu et al. have investigated that high-density stacking faults of improve dielectric permittivity and loss tangent of SiC nanowires.^{13−15} Wang et al. have synthesized SiC−Fe₃O₄

hybrid nanowires offering a much improved capability for microwave absorption.¹⁶ Surface modification and doping are also feasible methods to enhance the microwave absorption.17,18 Because of [i](#page-4-0)ts thermal stability and dielectric properties, SiC becomes a promising candidate as elevated tem[perat](#page-4-0)ure absorber. However, improving its dielectric properties and absorption capacity are still a great challenge.

In this letter, we successfully fabricated a novel hierarchical architecture of NiO nanorings on SiC. We report our original observation on dielectric properties in the temperature range from 373 to 773 K and X band. The results demonstrate that the microwave absorption is significantly superior to SiC, and the intensity and bandwidth are obviously enhanced. The mechanism via multirelaxation to enhance microwave absorption is also discussed.

The hierarchical architecture, NiO nanoring on SiC (NiO-SiC) in this work was fabricated by chemical deposition and

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Figure 1. Morphology of SiC and NiO-SiC and SAED pattern taken from NiO nanocrystals and SiC. (a) SEM image of SiC; (b) SEM image of NiO-SiC; (c−f) TEM images of NiO-SiC; (g) HRTEM image of NiO-SiC; (h) SAED pattern of NiO-SiC.

Figure 2. Complex permittivity and loss tangent versus frequency of (a−c) SiC and (d−f) NiO-SiC.

oxidization. The detailed experimental procedure is depicted in the Supporting Information. Figure 1 is the scanning electron microscopy (SEM) images of SiC and NiO-SiC, the transmis[sion electron microscopy](#page-3-0) (TEM) images of as-synthesized morphology, and selected area electron diffraction (SAED) pattern of NiO-SiC. From the SEM image in Figure 1a, SiC are of irregular shape with the size ranging from several hundred nanometers to several micrometers. The surfaces of SiC particles are fully covered by NiO shown in Figure 1b. From the TEM images (Figure 1c, d), it can be noted that NiO uniformly covers the surface of SiC, exhibiting a ringlike shape, with almost the same sizes. Figure 1d almost shows a whole assynthesized SiC particle, from which it can be seen that the particle is completely covered with NiO. Figure 1e, f show that some NiO nanorings aggregate together, which may be

correlated with activated spots produced in the reaction process. The external diameter of NiO nanorings is about 40−50 nm, and the internal diameter is approximately 20 nm. The clear lattice fringes of NiO nanorings illustrate that the interplanar spacing is about 0.21 nm, corresponding to (200) planes of face-centered cubic NiO, shown in Figure 1g. The lattice fringes indicate that the NiO nanorings are composed of highly crystallized nanocrystals. The SAED pattern further confirms that the specimen includes two phases: nickel oxide (blue circles) and silicon carbide (red circles), shown in Figure 1h. The structure of as-fabricated NiO-SiC, is identified by XRD patterns (Figure S1 in the Supporting Information). The XRD patterns and EDS spectra of SiC and NiO-SiC clearly demonstrate that there are no i[mpurity phases \(Figure S](#page-3-0)1 and S2 in the Supporting Information). From the EDS spectra

shown in Figure S2 in the Supporting Information, the atom percent of Ni and SiC is obtained (Table S1 in the Supporting Information), and the ma[ss percent of NiO is](#page-3-0) estimated between 1 and 3%. The actual percent of NiO on Si[C should be](#page-3-0) [much lower](#page-3-0) than the percent shown in the EDS spectra because the depth of EDS is far smaller than the size of SiC particles.¹⁵

The dielectric properties of NiO-SiC versus frequency at elevated temperatures are shown in Figure 2. The compl[ex](#page-4-0) permittivity ε and the loss tangent tg δ of SiC change slightly with the increase of temperature (Figure 2a−c[\).](#page-1-0) Meanwhile, the complex permittivity ε and the loss tangent tg δ of NiO-SiC demonstrate strong temperature depe[nd](#page-1-0)ence (Figure 2d−f). The real part ε' of NiO-SiC increases ∼90% and the imaginary part ε'' increases dra[m](#page-1-0)atically, more than four times from 373 to 773 K. The tgδ of NiO-SiC increases from 0.21 to 0.59 with the increase in temperature. Compared with the dielectric properties of SiC, the ε'' of NiO-SiC is almost three times higher than that of SiC and the maximum tg δ of NiO-SiC is more than four times higher than that of SiC at 773 K. The ε " of SiC and NiO-SiC indicates the evidence of relaxation.

According to Debye theory, the imaginary part ε'' could be regarded as the complex contributions of both relaxation and conductance, described as the following⁵

$$
\varepsilon''(\omega) = \varepsilon_p'' + \varepsilon_c'' = (\varepsilon_s - \varepsilon_\infty) \frac{\omega \tau}{1 + \omega^2 \tau^2} + \frac{\sigma(T)}{\varepsilon_0 \omega} \tag{1}
$$

where $\sigma(T)$ is temperature-dependent electrical conductivity, ω the angular frequency, τ the temperature-dependent relaxation time, ε _s the static permittivity and ε_{∞} the relative permittivity.

Consequently, ε'' contains two sections: ε''_p induced by relaxation and e_c'' produced by conductance.⁶ Figure 3a and 3b is the contribution of the ε_p'' on the ε'' of SiC and NiO-SiC at differe[nt](#page-4-0) temperatures, respectively. The contribution of the $\varepsilon_c^{\prime\prime}$ on the ε'' is illustrated in the Supporting Information (Figure S3).

As shown in Figure 3a, th[ere are two relaxations](#page-3-0) of SiC, Relaxation I and Relaxation II, which are located at the frequencies of 8.9 and 10.9 GHz, respectively. The relaxation of SiC mainly originated from the defect dipole polarization and the interfacial polarization.^{5,20} The defect dipoles illustrated in Figure 4a are generated by the charge unbalance around the carbon vacancies in the Si[C l](#page-4-0)attice. The interfacial dipoles are also re[su](#page-3-0)lted from the charge unbalance, which exists among the SiC interfaces. The defect polarizations and the interfacial polarizations are coexist in the SiC shown in Figure 4c. According to the previous literature, the relaxation time τ for grain boundary regions is much larger than that for bulks.²¹ [Th](#page-3-0)e grain boundary response relaxes at lower frequencies than the bulk. Therefore, Relaxation I is induced by the int[erf](#page-4-0)acial polarization between SiC, and Relaxation II is induced by the defect polarization in the SiC (Figure 4a and 4c).

Figure 3b demonstrates multirelaxations in NiO-SiC, defined as Relaxation III, Relaxation IV, Relax[at](#page-3-0)ion V[,](#page-3-0) and Relaxation VI at the frequencies of 8.5, 9.9, 10.9, and 11.8 GHz, respectively. Multirelaxations originate from (1) the defect polarization in NiO nanocrystal, $2^{2}(2)$ the defect polarization in SiC grains^{20} and (3) the multiple interfacial polarization in NiO-SiC.^{23,24} These polarizatio[ns](#page-4-0) are depicted in Figure 4b, d.

The rela[xat](#page-4-0)ion time τ has the relationship with temperature and activ[ation](#page-4-0) energy, expressed as the following²

Figure 3. ε_p'' versus frequency of (a) SiC and (b) NiO-SiC at different temperatures.

where τ_0 is the relaxation time at infinite temperature, U the activation energy for relaxation, T the absolute temperature, and k the Boltzmann constant.

Although Relaxation V derives from the defects polarization of SiC because it has the same frequency as Relaxation II in SiC, Relaxation VI derives from the relaxation of NiO nanocrystal defects. According to eq 2, the activation energy U of NiO is slightly smaller than that of the SiC, 26,27 so the relaxation frequency of NiO defects will be higher than that of SiC. Relaxations III and IV are induced by th[e m](#page-4-0)ultipleinterfacial polarization in NiO-SiC (Figure 4b, d). δ

The reflection loss R_L of SiC and NiO-SiC is shown in Figure 5a and the Supporting Information (Figur[e](#page-3-0) S4) [in](#page-4-0) the range from 373 to 773 K and X band. The minimum R_L and the

Figure 4. Schematic illustrations of dipole polarization in SiC and NiO-SiC. (a) Defect dipole polarization of SiC; (b) defect dipole polarization of NiO-SiC; (c) multipolarization of SiC; (d) multipolarization of NiO-SiC.

maximum bandwidth of SiC are approximately −15 dB and ∼2.5 GHz at the level of approximately −10 dB (Figure 5b), respectively, whereas the bandwidth is zero at the level of approximately −20 dB in the temperature range (Figure 5c). The as-fabricated NiO-SiC exhibits superior microwave absorption in X band at elevated temperature. The minimum R_L of NiO-SiC is approximately -50 dB, more than three times than that of SiC. The bandwidth of NiO-SiC covers the whole X band at 673 and 773 K at the level of approximately −10 dB (Figure 5b), and reaches 3.6 GHz at the level of approximately −20 dB (Figure 5c). The bandwidth of NiO-SiC exhibits a remarkable widening tendency with the increase in temperature.

For comparison, the tg δ and reflection loss R_L of NiO/SiC mixture and NiO-SiC architecture are shown in the Supporting Information (Figures S5 and S6) in X band. The tg δ of NiO-SiC architecture is much higher than that of NiO/SiC mixtures. The reflection loss R_L of NiO-SiC is superior to that of the mixtures. The microwave absorption of NiO-SiC architecture results from NiO-SiC hierarchical architecture, as well as NiO nanorings.²⁸

The microwave absorption of NiO-SiC is attributed to the synergistic [e](#page-4-0)ffect between multirelaxation and conductance as shown in Figure 5d, e (details shown in Figures S7 and S8 in the Supporting Information). The tg δ_{ep} and tg δ_{ec} are loss tangent values of contribution from relaxation and conductance, respectively. The tg δ_{ep} is the dielectric loss induced by relaxation, determined by ε_p'' with multirelaxation (Figure 3). The tg δ_{ec} is the dielectric loss produced by conductance. The tg $\delta_{\rm{ep}}$ and tg $\delta_{\rm{ec}}$ evidently rise with the increase in temperat[ure](#page-2-0). With the increase in frequency, the tg δ_{ep} ascends and the tg δ_{ec} descends. The tg $\delta_{\rm{ep}}$ makes an important contribution to the tgδ. The multirelaxation requires the expenditure of microwave

Figure 5. (a) Reflection loss R_1 of SiC and NiO-SiC versus frequency $(T = 673 \text{ K})$; absorption bandwidth of SiC and NiO-SiC at different temperatures at the level of (b) -10 dB and (c) -20 dB; (d) loss tangent tg $\delta_{\rm ec}$ and tg $\delta_{\rm ep}$ versus frequency (T = 673 K); (e) loss tangent tg $\delta_{\rm ec}$ and tg $\delta_{\rm ep}$ versus temperature ($f = 11.2$ GHz).

energy at different frequencies.²⁹ Therefore, the bandwidth and intensity of microwave absorption are enhanced by multirelaxation and conductance. [Th](#page-4-0)e synergistic effect between relaxation and conductance tunes the microwave absorption of NiO-SiC.

In summary, NiO hierarchical nanorings on SiC are successfully fabricated. The nanorings obviously enhance the dielectric properties of SiC. The imaginary part ε'' and loss tangent tg δ dramatically increases from 2.2 to 10.0 and from 0.21 to 0.59 with the increasing temperature, respectively. The NiO-SiC architecture exhibits superior microwave absorption with both high efficiency and broad frequency at elevated temperature. The bandwidth of NiO-SiC is significantly wider than that of SiC. The minimum R_L reaches to approximately −50 dB. The tuned absorption arouses from the synergistic effect between multirelaxation and conductance. These findings open a practicable pathway to tune microwave absorption.

ASSOCIATED CONTENT

6 Supporting Information

Experimental details, XRD and EDS spectra, curves of $\varepsilon_{\rm c}''$ versus frequency, the reflection loss of SiC and NiO-SiC, the loss tangent tg $\delta_{\rm ec}$ and tg $\delta_{\rm ep}$ of NiO-SiC. This material is available free of charge via the Internet at http://pubs.acs.org/.

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Notes

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